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Accelerating beam generation by dielectric metasurfaces fabricated through low-temperature ALD

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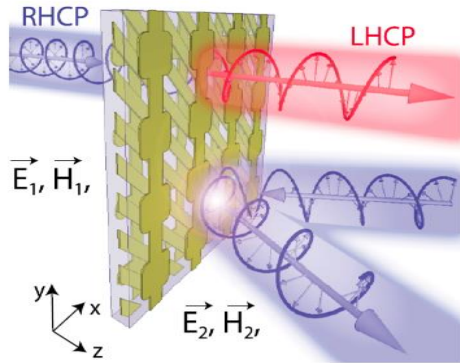
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Outline

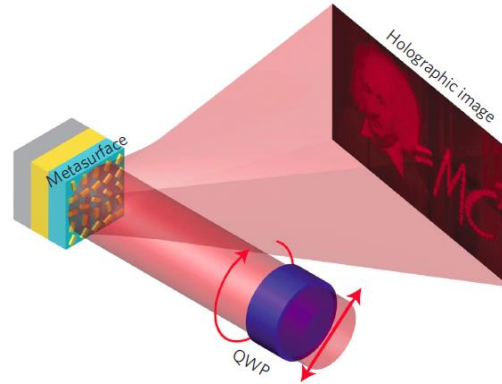
- ◀ Background & Motivation for Optical Metasurface
- ◀ All Dielectric Metasurface Device design
- ◀ All Dielectric Metasurface Device fabrication
- ◀ Optical Characterization
- ◀ Conclusion

Applications of Optical Metasurfaces



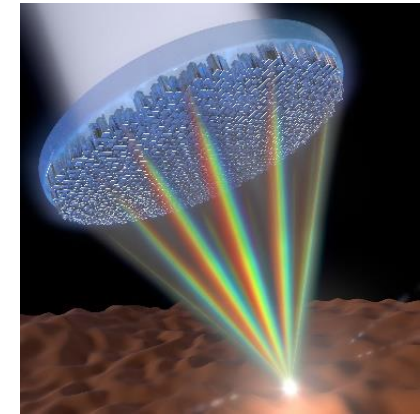
Metallic structure for Light Asymmetric Transmission

Pfeiffer, *et al.*, **Phys. Rev. Lett.**, 2014



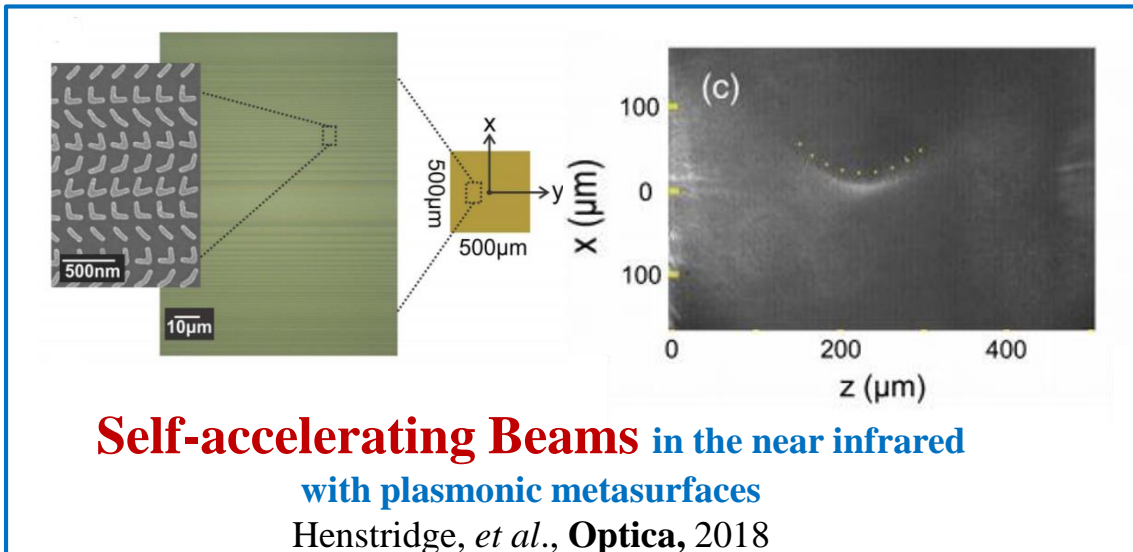
Metallic structure on metal mirror for Holographic display

Zheng, *et al.*, **Nat Nanotechnol.**, 2015



Dielectric on GaN for Optical Imaging

Chen, *et al.*, **Nat Nanotechnol.**, 2018



Self-accelerating Beams in the near infrared with plasmonic metasurfaces

Henstridge, *et al.*, **Optica**, 2018

Self-accelerating beams

- An exotic **diffraction-free** solution to Maxwell's equations
- A wave packet propagating along a **curved trajectory**
- **This work: Spin-controlled generation of self-accelerating beams in the visible with all dielectric metasurfaces**

All-dielectric Metasurface for Self-accelerating Beam

◀ Motivation

- It is a novel idea to apply all-dielectric metasurface for self-accelerating beam at visible light range
- Compact wafer level generation of self-accelerating beam compared to traditional methods; semiconductor process flow can be applied in device fabrication

◀ Benefits of all-dielectric metasurface

- Reduced optical loss compared to plasmonic structure
- Be operated under transmission mode

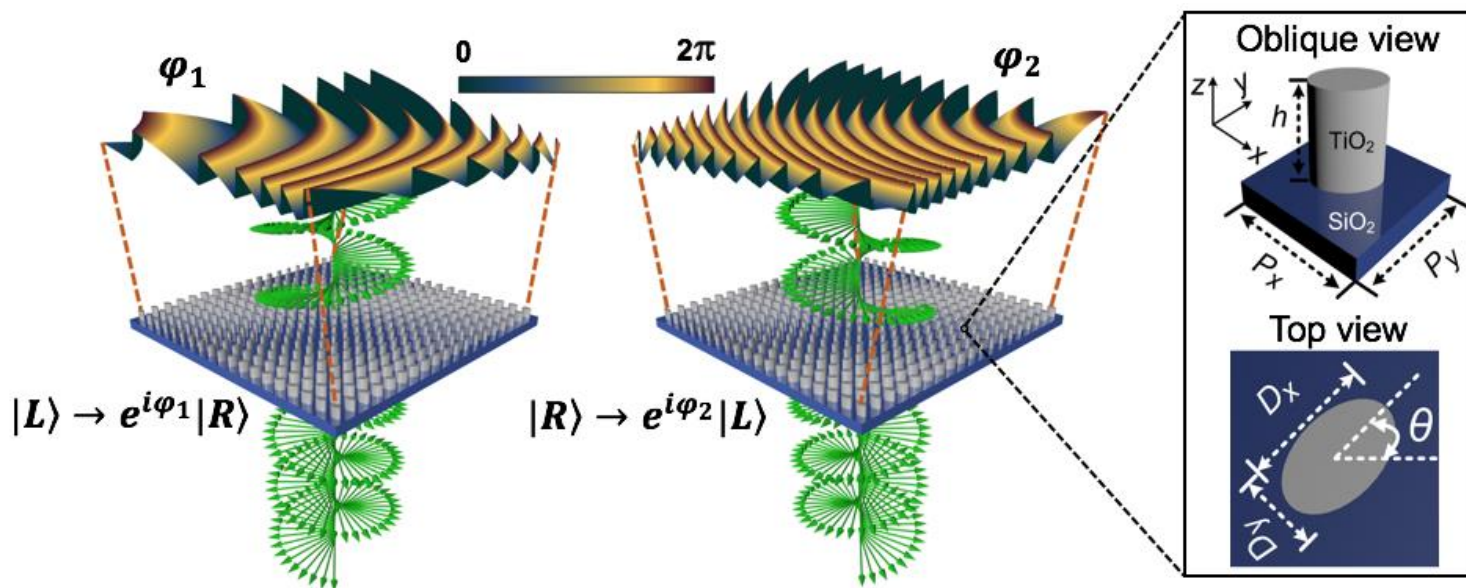
Self-accelerating beam generation with a phase mask

Accelerating beams propagating along a certain trajectory can be generated with associated phase masks

Acceleration profile	Applied phase
Parabolic: $c(z) = az^2$	$\phi(y) = -4/3 a^{1/2} k y^{3/2}$
Quartic: $c(z) = az^4$	$\phi(y) = -16/21 (3a)^{1/4} k y^{7/4}$
Logarithmic: $c(z) = a \ln(bz)$	$\phi(y) = e^{-1} a^2 b k (1 - \exp[-y/a])$
Polynomial: $c(z) = az^n$ (for n even)	$\phi(y) = kn^2 y^2 \frac{[a(1-n)/y]^{1/n}}{(2n-1)(1-n)}$

L. Froehly, *et al.*, *Arbitrary accelerating micron-scale caustic beams in two and three dimensions*, **Optics Express**, 2011

Spin-controlled self-accelerating beam generation



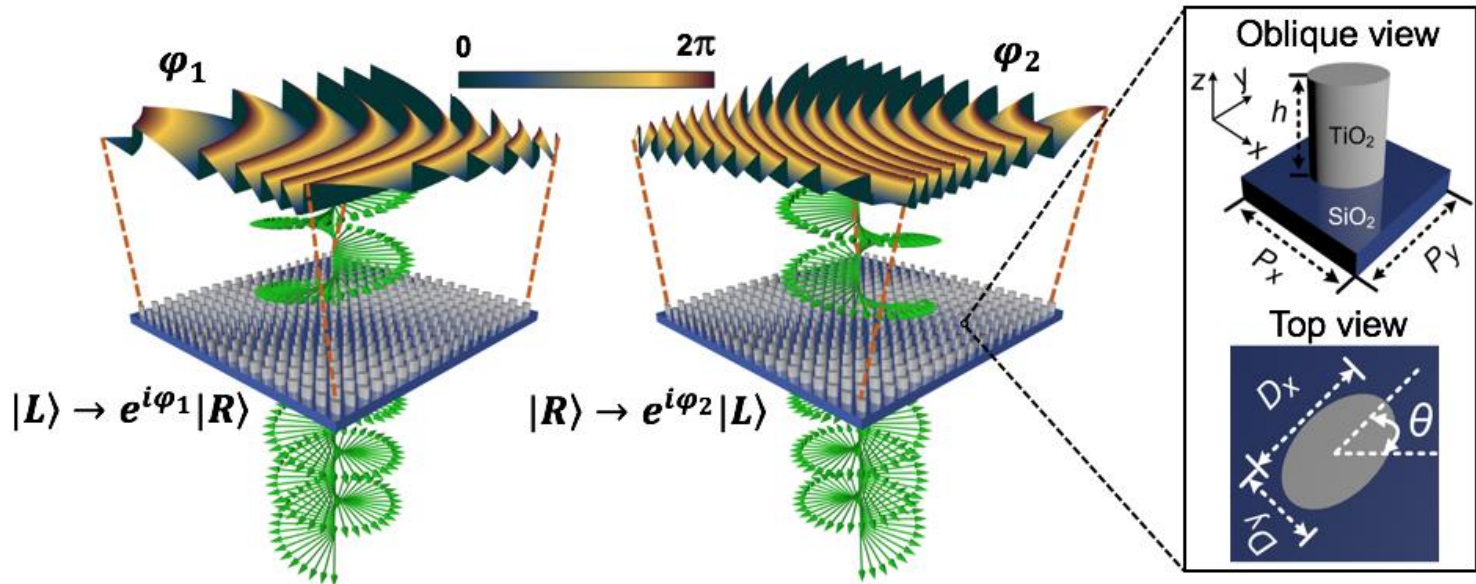
Target: design an all dielectric plate to independently output phase profiles $\varphi_1(x, y)$ and $\varphi_2(x, y)$ for the input LCP and RCP light

The light transformation characteristics of such a metasurface can be described by a Jones Matrix $J(x, y)$

$$e^{i\varphi_1(x,y)}|\mathbf{R}\rangle = J(x, y)|\mathbf{L}\rangle$$

$$e^{i\varphi_2(x,y)}|\mathbf{L}\rangle = J(x, y)|\mathbf{R}\rangle$$

Spin-controlled self-accelerating beam generation



The Jones Matrix can be further expressed as:

$$J(x, y) = \frac{1}{2} \begin{bmatrix} e^{i\varphi_1(x,y)} + e^{i\varphi_2(x,y)} & ie^{i\varphi_2(x,y)} - ie^{i\varphi_1(x,y)} \\ ie^{i\varphi_2(x,y)} - ie^{i\varphi_1(x,y)} & -e^{i\varphi_1(x,y)} - e^{i\varphi_2(x,y)} \end{bmatrix}$$



Requiring the nanostructures to have:

$$\begin{cases} \delta_x(x, y) = [\varphi_1(x, y) + \varphi_2(x, y)]/2 \\ \delta_y(x, y) = [\varphi_1(x, y) + \varphi_2(x, y)]/2 - \pi \\ \theta(x, y) = [\varphi_1(x, y) - \varphi_2(x, y)]/4 \end{cases}$$

Design of elliptical pillars

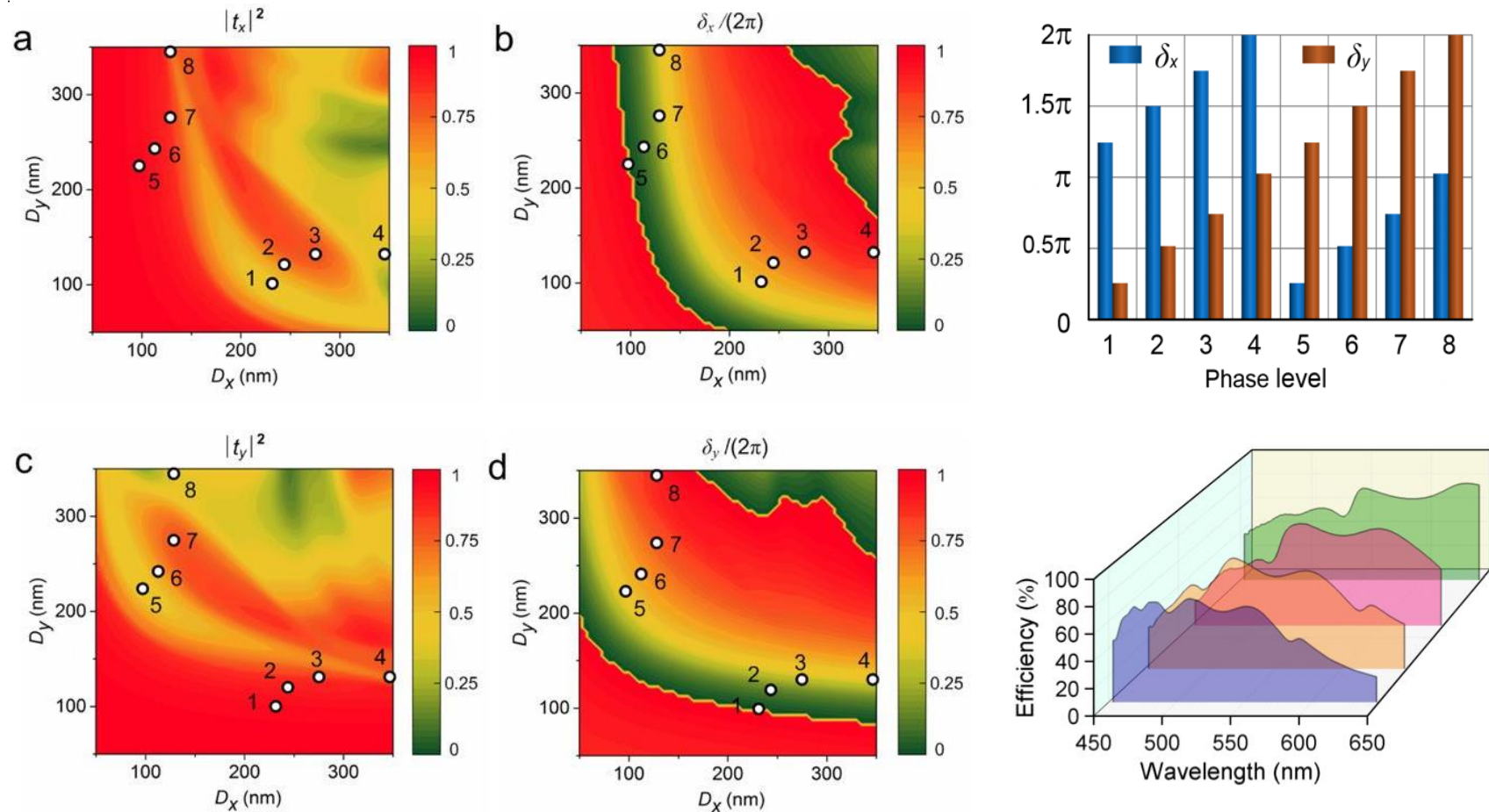
◀ Choice of Material

- High refractive index in visible light range (TiO₂: 2.61)
- Low adsorption in the visible light range (TiO₂: 0)
- Amorphous material is preferred

◀ Choice of Nano Structure

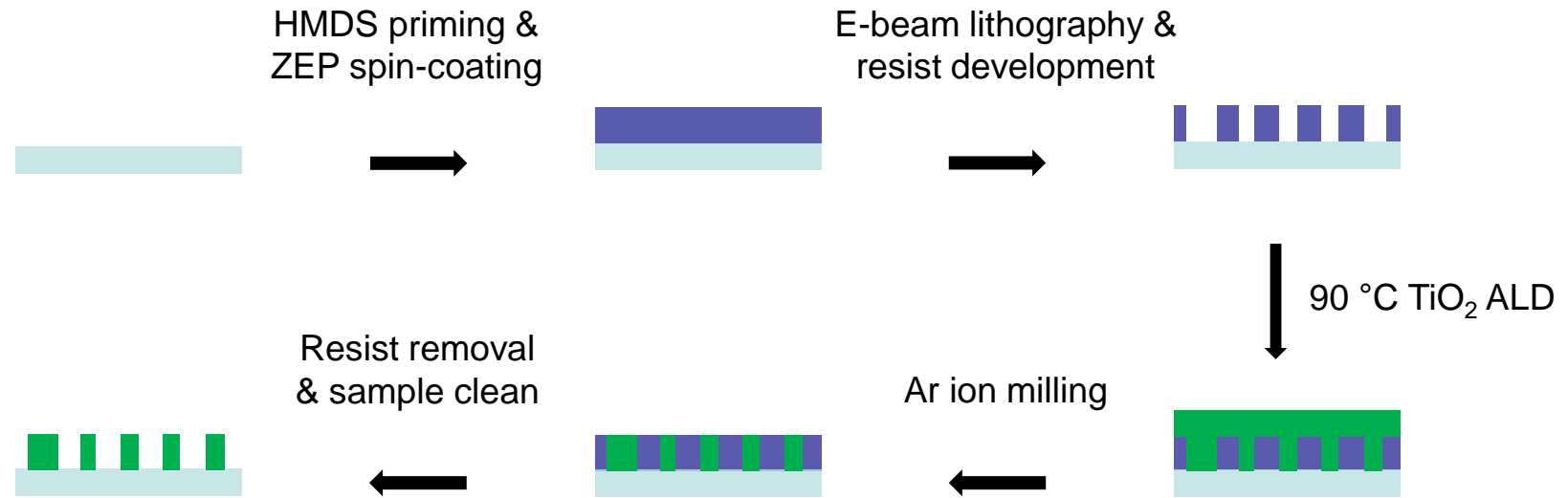
- Nano pillars with defined Dx, Dy and orientation angle which can be easily created by E-beam litho.
- Z 600nm with choice of high refractive index material so it is not high AR structure

Design of elliptical pillars



- Four fundamental TiO₂ pillars & their rotated structures along the y-axis
- All pillars acts closely as a half-wave plates with the required phase shift along the two principle axes

Device Fabrication



❑ Requirement of the TiO₂ ALD

- Low-temperature operation (no resist reflow)
- Plasma-free operation (no resist deterioration)
- AR: no more than 6:1 by design (z 600nm while Dx or Dy in between 100nm and 350nm)

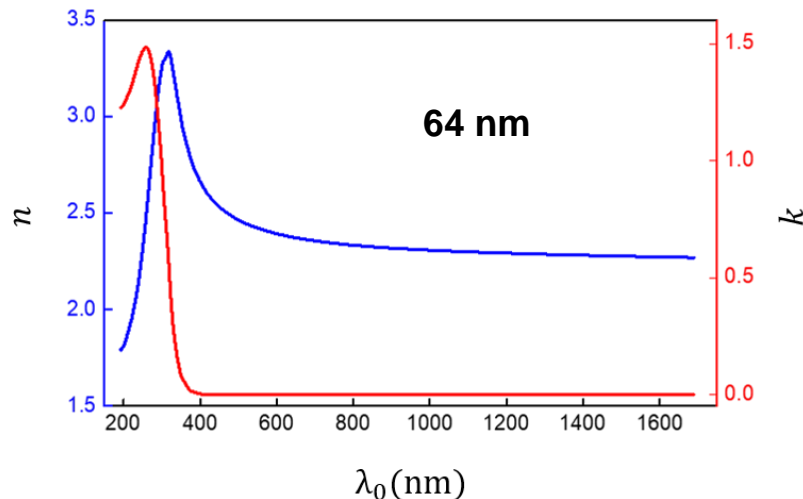
Low-temperature TiO₂ ALD

□ Deposition:

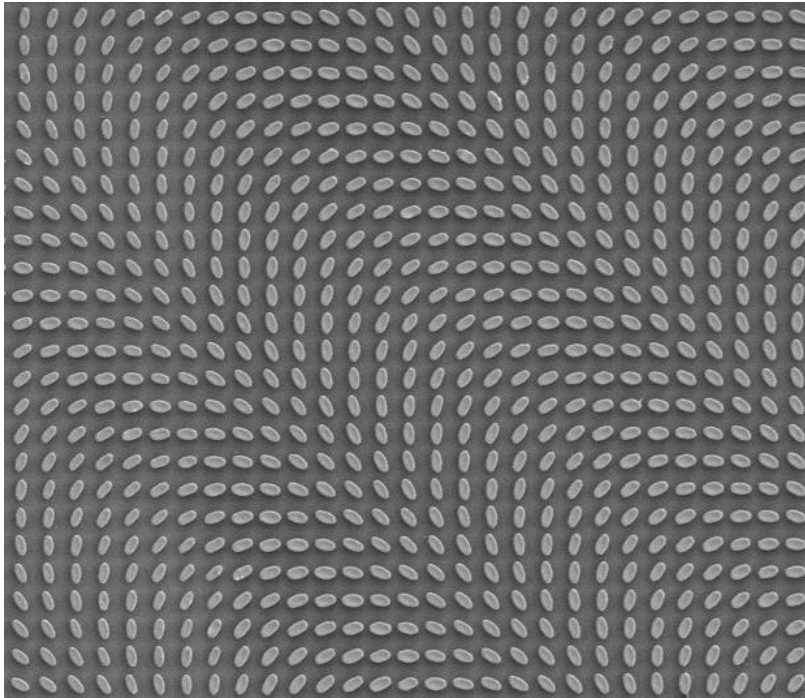
- Arradiance GEMStar™
- Fused silica wafers and Si witnesses
- 200nm of TiO₂ by TDMATi with water (90C)

□ Measurement:

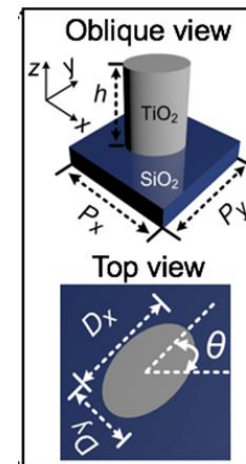
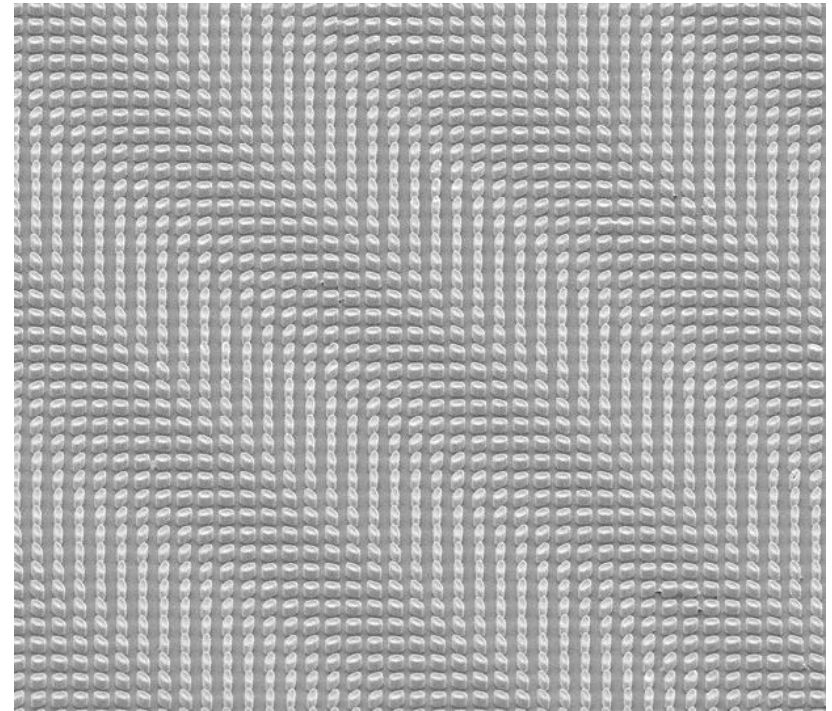
- Thickness by ellipsometer
- Optical constants by ellipsometer
- Imaging by SEM



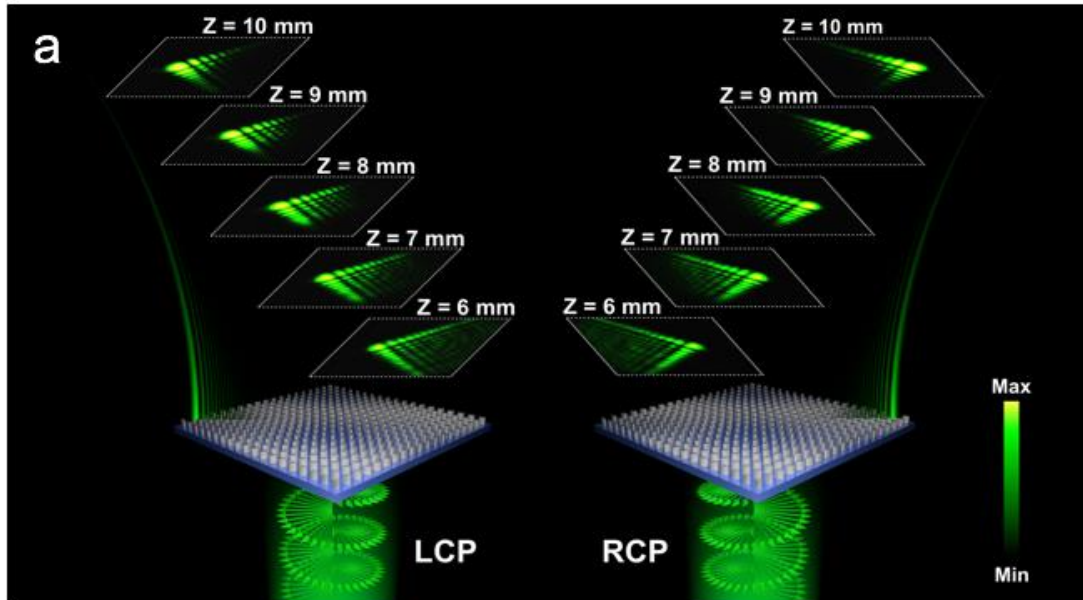
Fabricated TiO₂ Metasurfaces



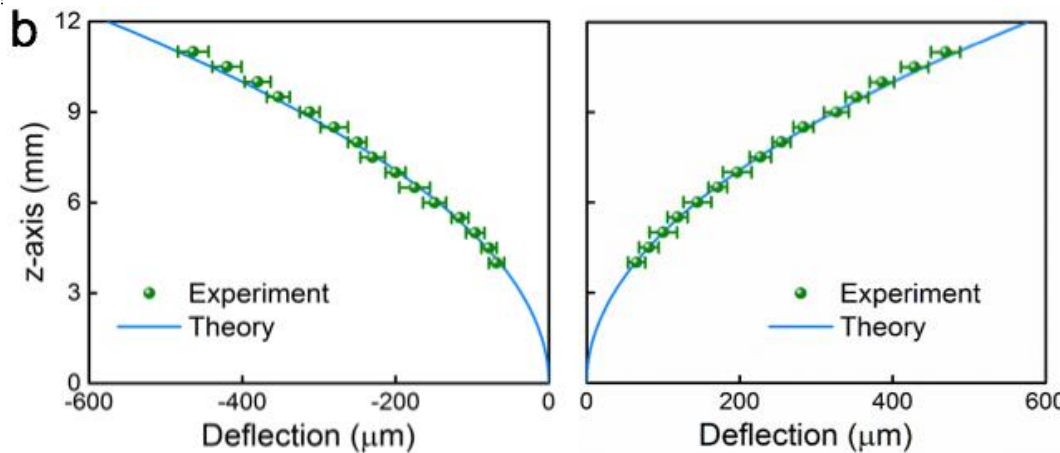
450nm



Example 1: spin-controlled propagation



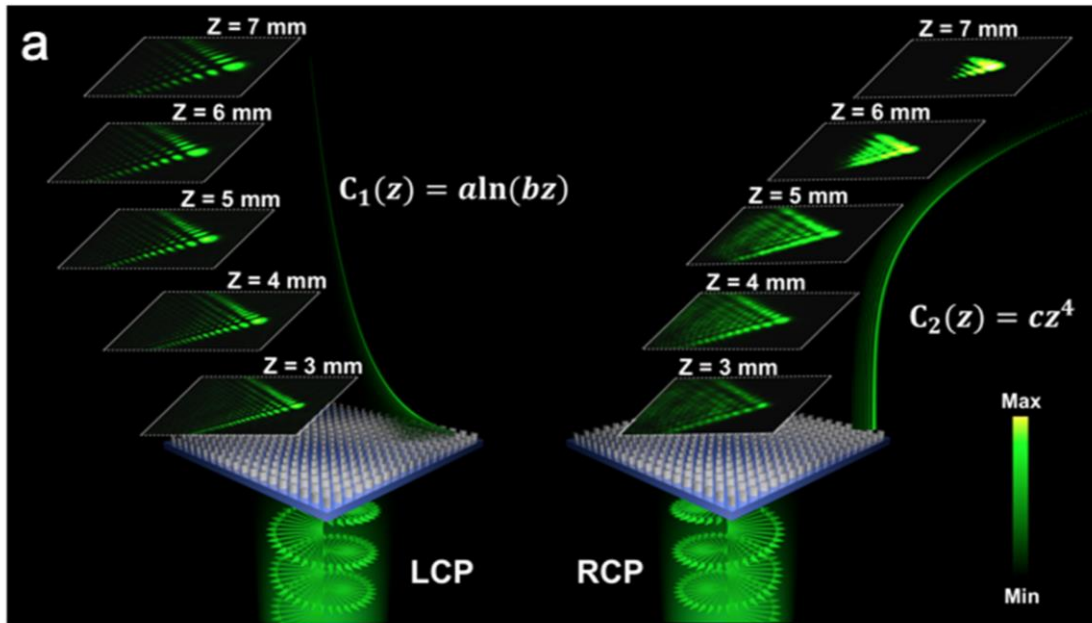
$$\begin{cases} x = nZ^2 \\ y = nZ^2 \end{cases} \quad n = 4 \text{ m}^{-1}$$



$$\varphi_1(x, y) = -\frac{4}{3} n^{\frac{1}{2}} k (x^{\frac{3}{2}} + y^{\frac{3}{2}})$$

$$\varphi_2(x, y) = \varphi_1(w - x, w - y)$$

Example 2: spin-controlled beam switching



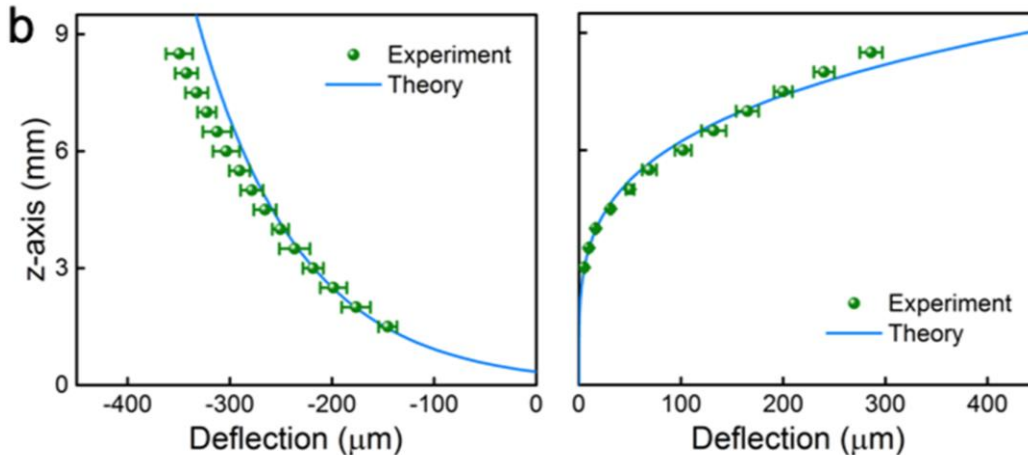
$$C_1: \begin{cases} x = a \ln(bz) \\ y = a \ln(bz) \end{cases}$$

$$C_2: \begin{cases} x = cz^4 \\ y = cz^4 \end{cases}$$

$$a = -10^{-4} \text{ m}$$

$$b = 8 \times 10^3 \text{ m}^{-1}$$

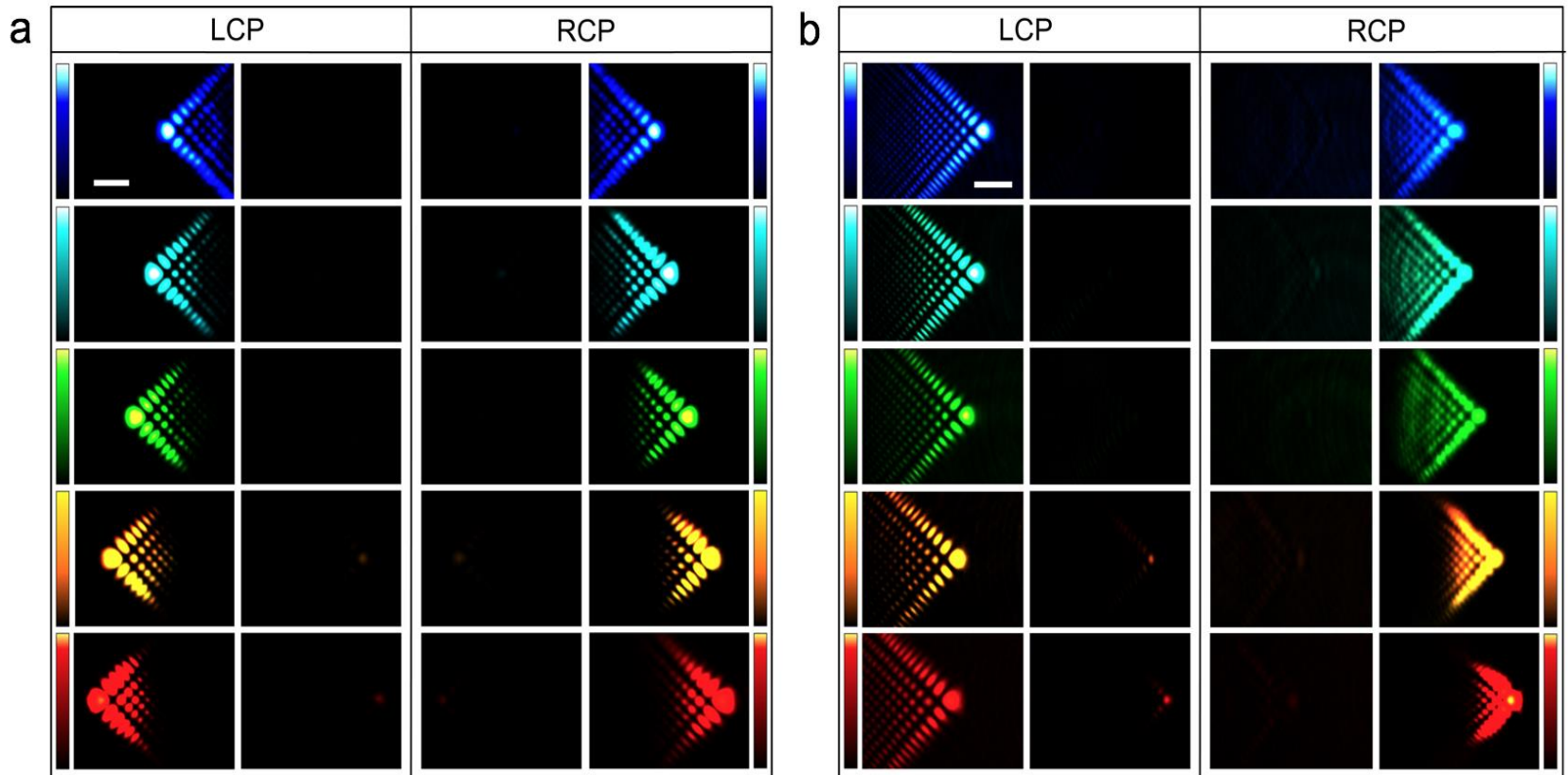
$$c = 6.7 \times 10^4 \text{ m}^{-3}$$



$$\varphi_1(x, y) = e^{-1} a^2 b k [2 - \exp(-\frac{x}{a}) - \exp(-\frac{y}{a})]$$

$$\varphi_2(x, y) = -\frac{16}{21} (3c)^{\frac{1}{4}} k (x^{\frac{7}{4}} + y^{\frac{7}{4}})$$

Example 3: broadband operation in the visible



Summary

- High efficiency and versatile generation of self-accelerating beams with a TiO_2 -based all-dielectric metasurface.
- By changing the design of metasurface, broadband and spin controlled self-accelerating beams can be achieved.
- Successful device fabrication through a semiconductor process incorporating low-temperature ALD.

Further reading: Fan, et al., Broadband generation of photonic spin-controlled arbitrary accelerating light beams in the visible, *Nano Letters*, 2019